

Meteotsunamis and Proudman resonance

MELINAND Benjamin
IMB, Université de Bordeaux

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 - b) The Proudman resonance
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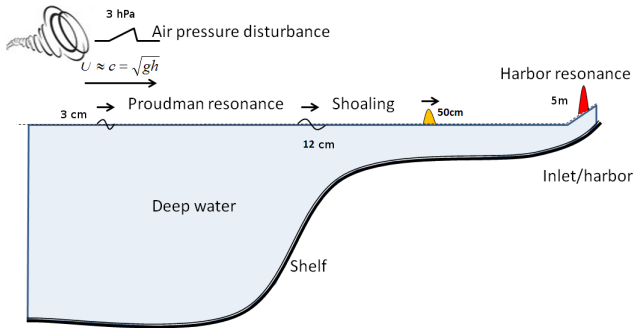
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- Principal causes : earthquakes, volcanic eruptions, underwater explosions.
- Typical wavelength : 100 kilometers.
- Amplitude offshore : 50 centimeters.

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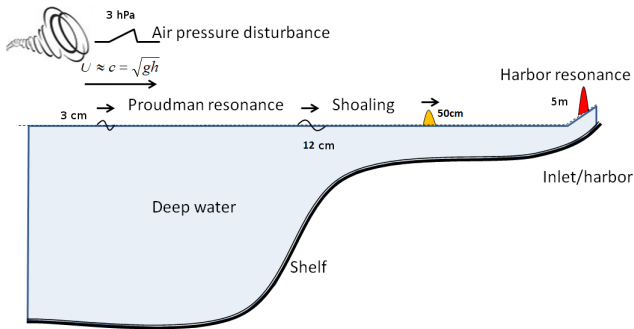
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Meteotsunami



One example : Nagasaki(1979), "Abiki", pressure disturbance traveled at 22-39 m/s. Amplitude on the coast : 5m.

Mechanisms

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- We consider a disturbance which traveling at the speed U .
- The typical speed of long ocean waves $c = \sqrt{gH}$ where H is the typical depth.
- There is a resonance if U is close to c .

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a) Mathematical formulation

b) Non dimensionalized equations

Assumptions and physical law

- We consider a fluid occupying a domain Ω_t in \mathbb{R}^{d+1} limited from below by a fixed bottom and above by a free surface.

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- The fluid particles do not cross the bottom or the surface.
- At the surface, we have a pressure $P_0 = P_0(x, t)$.
- The water depth is bounded from below by $h_{\min} > 0$.

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The Zakharov/Craig-Sulem formulation

- We consider the following standard problem :

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- If $\zeta = 0$ and $b = 0$, $\widehat{G[\zeta, b](\psi)}(\xi) = |\xi| \tanh(H|\xi|) \widehat{\psi}(\xi)$.

We can reformulate the whole problem with the two unknown variables (ζ, ψ) :

$$\begin{cases} \partial_t \zeta - G[\zeta, b](\psi) = 0 \\ \partial_t \psi + \zeta + \frac{1}{2} |\nabla_x \psi|^2 - \frac{(G[\zeta, b](\psi) + \nabla_x \zeta \cdot \nabla_x \psi)^2}{2(1 + |\nabla_x \zeta|^2)} = -\frac{P_0}{\rho}. \end{cases}$$

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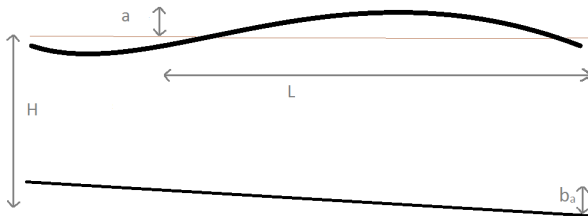
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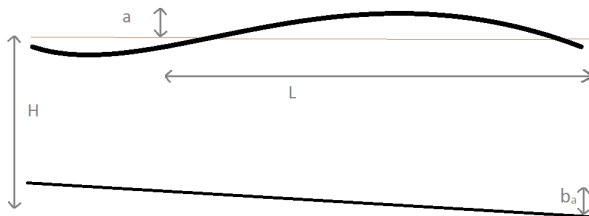
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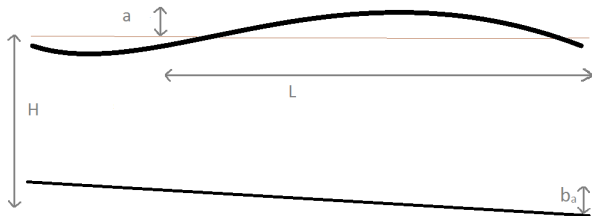
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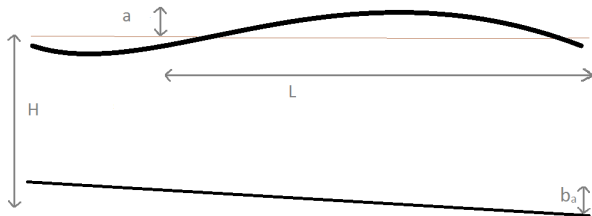
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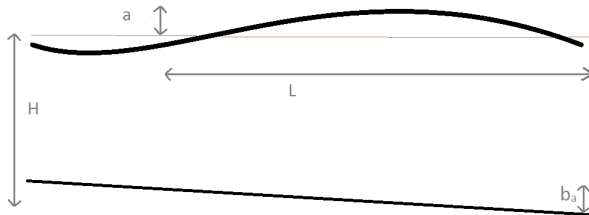
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- The bathymetric parameter : $\beta = \frac{b_a}{H}$.
- The typical speed : $c = \sqrt{gH}$ (in the shallow-water case).

We can renormalize the variables :

$$x' = \frac{x}{L}, \quad z' = \frac{z}{H}, \quad \zeta' = \frac{\zeta}{a}, \quad \dots$$

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In the case of the Proudman resonance :

- $L \sim 30\text{-}100\text{km}$, $H \sim 50\text{-}500\text{m}$ and $a \sim 3\text{-}20\text{cm}$.

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- $L \sim 30\text{-}100\text{km}$, $H \sim 50\text{-}500\text{m}$ and $a \sim 3\text{-}20\text{cm}$.
- $\varepsilon \sim 10^{-3}$, $\mu \sim 10^{-5}\text{-}10^{-6}$, $c \sim 22\text{-}70\text{m/s}$.

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Rayleigh-taylor criterion

We can compute the pressure P in Ω_t . We denote by $\alpha(\zeta, \psi) := -(\partial_z P)|_{z=\varepsilon\zeta}$, called the Rayleigh-Taylor coefficient.

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If we linearize the water waves equations around (ζ_0, ψ_0) , we have

$$\partial_t \begin{pmatrix} \zeta \\ \psi \end{pmatrix} + \begin{pmatrix} \varepsilon \underline{V}_0 \cdot \nabla & \frac{1}{\mu} G_\mu[\varepsilon \zeta_0, \beta b](\cdot) \\ \alpha(\zeta_0, \psi_0) & \varepsilon \underline{V}_0 \cdot \nabla \end{pmatrix} \begin{pmatrix} \zeta \\ \psi \end{pmatrix} + \text{lower term} = \begin{pmatrix} 0 \\ -P \end{pmatrix}$$

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If we denote by $A := \begin{pmatrix} i\varepsilon \underline{V}_0 \cdot \xi & \frac{1}{\mu} G_\mu[\widehat{\varepsilon \zeta_0}, \beta b] \\ \alpha(\zeta_0, \psi_0) & i\varepsilon \underline{V}_0 \cdot \xi \end{pmatrix}$, we need that $Sp(A) \subset i\mathbb{R}$ (hyperbolicity).

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Therefore, we need that $\alpha(\zeta_0, \psi_0) > 0$.

Existence and uniqueness of the Water waves equation

For $s \geq 0$, we define $\dot{H}^s(\mathbb{R}^d) := \{u \in L^2_{\text{loc}}(\mathbb{R}^d), \nabla u \in H^{s-1}(\mathbb{R}^d)\}$.

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Theorem

Let N large enough.

Let $P \in \mathcal{C}^0(\mathbb{R}^+, H^N(\mathbb{R}^d))$, $(\psi_0, \zeta_0) \in \dot{H}^N(\mathbb{R}^d) \times H^N(\mathbb{R}^d)$, such that

$$1 + \varepsilon \zeta_0 + \beta b_0 \geq h_{\min} \text{ and } \underline{\alpha}(\psi_0, \zeta_0) \geq \alpha_0.$$

There exists $T > 0$, $(\psi, \zeta) \in \mathcal{C}^0\left([0, \frac{T}{\max(\varepsilon, \beta)}], \dot{H}^N(\mathbb{R}^d) \times H^N(\mathbb{R}^d)\right)$ unique solution of the water waves equation with the initial data (ψ_0, ζ_0) .

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An asymptotic model

We suppose that ε, β, μ are small. We have

$$\left\{ \begin{array}{l} \partial_t \zeta - \overbrace{\frac{1}{\mu} G[\varepsilon \zeta, \beta b](\psi)}^{\sim -\partial_x^2 \psi} = 0 \\ \partial_t \psi + \zeta + \underbrace{\frac{\varepsilon}{2} |\nabla_x \psi|^2 - \varepsilon(\dots)}_{\sim \varepsilon C(|\zeta|_{HN}, |b|_{HN}) \|\psi\|_{HN}} = -P_0. \end{array} \right.$$

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Proposition

For N large enough, there exists $T > 0$ and a solution of the water waves equation $(\psi, \zeta) \in \mathcal{C}^0\left([0, \frac{T}{\max(\varepsilon, \beta)}\right], \dot{H}^{N+1}(\mathbb{R}^d) \times H^N(\mathbb{R}^d))$,

$$\left\{ \begin{array}{l} \partial_t \zeta + \partial_x^2 \psi = \max(\varepsilon, \beta, \mu) R, \\ \partial_t \psi + \zeta = -P_0 + \max(\varepsilon, \beta, \mu) S. \end{array} \right.$$

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Proudman resonance

We consider the following equation

$$\begin{cases} \partial_t \zeta + \partial_x^2 \psi = 0, \\ \partial_t \psi + \zeta = -P_0, \\ \zeta|_{t=0} = 0, \\ \psi|_{t=0} = 0, \end{cases}$$

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We have the Proudman resonance.

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- And the landslide-generated tsunamis ?